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Journal of the European Ceramic Society 28 (2008) 1899–1907

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# Arc-jet testing on  $HfB_2$  and  $HfC$ -based ultra-high temperature ceramic materials

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Received 25 July 2007; received in revised form 13 November 2007; accepted 25 November 2007 Available online 4 March 2008

# **Abstract**

The behaviour of pressureless sintered HfC and HfB<sub>2</sub> ceramics, when exposed to high enthalpy plasma flows typical of atmospheric re-entry environment, was investigated with an arc-jet facility at temperatures exceeding 2000 ◦C. The surface temperature and emissivity of the materials were evaluated during the test. The microstructure modifications were analysed after exposure. Fluid dynamic numerical simulations were carried out to evaluate the catalytic atom recombination efficiencies of the materials at the experimental conditions. Surface and cross sections of the samples showed the formation of scales mainly consisting of HfO<sub>2</sub> and  $SiO_2$ . For the HfB<sub>2</sub>-based composite numerical results correlated quite well with experimental ones assuming a low catalytic surface behaviour. For the HfC-based material the surface behaviour changed from low catalytic to partially catalytic as the temperature increased. The post-test analyses confirm the potential of these composites to endure re-entry conditions with temperature approaching 2000 °C or even higher.

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Keywords: Arc-jet; Thermal protection systems; HfB<sub>2</sub>; HfC; Oxidation

# **1. Introduction**

Ultra-high temperature ceramics (UHTCs) are currently con-sidered as emerging materials for aerospace applications.<sup>[1–4](#page-7-0)</sup> The increasing attention is driven by the demand of developing reusable hot structures as thermal protection systems (TPS) of re-entry vehicles characterised by sharp leading edges and therefore by larger aerothermal heating than blunt edges, such as those on the Space Shuttle, able to withstand temperatures that may exceed 2000 ◦C during re-entry. As available materials cannot survive such extreme temperatures, new ones are required for advanced thermal protection systems.[1,4,5](#page-7-0) The use of UHTCs for sharp leading edges would also imply lower aerodynamic drag, improved flight performances and crew safety, due to the larger cross range and manoeuvrability along with more gentle re-entry trajectories.[3,6,7](#page-7-0)

Hafnium boride and hafnium carbide, belonging to the class of the UHTCs, are candidates for thermal protection materials

0955-2219/\$ – see front matter © 2008 Elsevier Ltd. All rights reserved. doi[:10.1016/j.jeurceramsoc.2007.11.021](dx.doi.org/10.1016/j.jeurceramsoc.2007.11.021)

in both re-entry and hypersonic vehicles because of their high melting points ( $\sim$ 3900 °C) and excellent chemical stability.<sup>[8–13](#page-7-0)</sup> Other notable properties are their high hardness, high electrical and thermal conductivity. $8-13$  Despite all the potentialities, so far these compounds have not been developed on industrial scale due to the difficult sinterability and low fracture toughness. Recent studies have pointed out that the addition of MoSi2 as sintering aid allows the achievement of highly dense bodies (98%) at 1950 °C by pressureless sintering.<sup>[14,15](#page-8-0)</sup> Furthermore, the addition of  $MoSi<sub>2</sub>$  is expected to improve the oxidation resistance due to the development of a silica protective coating[.16](#page-8-0)

In this paper, arc-jet testing at temperatures between 1950 ◦C and 2400  $\degree$ C is carried out on pressureless sintered HfB<sub>2</sub> and HfC-based materials. Arc-jet testing represents the best groundbased simulation of a re-entry environment, in different ways. On one hand, it provides the possibility to explore the oxidation behaviour of these materials under extreme conditions. On the other hand, the materials response to large heat fluxes is evaluated through the determination of two important parameters, i.e. emissivity and catalytic efficiency. High values of emissivity and low values of surface catalytic efficiency are desired for

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<span id="page-1-0"></span>





Fig. 1. Hemispheric HfB2 model used for arc-jet testing, curvature radius  $R = 7.5$  mm.

the above mentioned applications as they reduce temperature gradients and thermal stresses in the structure, thus enabling the vehicle to operate under relatively high enthalpy flow conditions. So far,  $ZrB_2-SiC$  and  $HfB_2-SiC$  composites were the predecessor materials analysed by arc-jet testing in the literature of UHTCs.<sup>[3,17](#page-7-0)</sup>

Microstructural modifications induced by high thermal loading are investigated and discussed. In addition, fluid dynamic numerical simulations are carried out in order to rebuild, through computational fluid dynamic (CFD) modelling, the experimental tests and to evaluate an average catalytic efficiency of the different materials with respect to oxygen and nitrogen surface recombination reactions.

### **2. Experimental**

# *2.1. Material processing and characterisation*

The following materials were selected for the arc-jet tests:

 $HfB_2 + 5$  vol% MoSi<sub>2</sub>, labelled as HB5.  $HfC + 5$  vol% MoSi<sub>2</sub>, labelled as HC5.

Commercial powders were used to prepare the ceramic materials: HfB<sub>2</sub> (Cerac Incorporated, Milwaukee, USA), particle size range  $0.5-5 \mu m$ , impurities: Al  $(0.07\%)$ , Fe  $(0.01\%)$ , Zr (0.47%); Cubic HfC (Cerac Inc., USA), 325 mesh, fisher size 1.1  $\mu$ m, grain size range 0.2–1.5  $\mu$ m; Tetragonal MoSi<sub>2</sub> (<2  $\mu$ m, Aldrich, USA), mean particle size  $2.8 \mu m$ , grain size range 0.3–5  $\mu$ m and oxygen content ∼1 wt%.

The powder mixtures were ultrasonically treated and milled for 24 h in absolute ethanol using zirconia milling media, then dried in a rotary evaporator and sieved to −250 mesh screen size. Four-centimetre diameter pellets were linearly pressed and subsequently cold isostatically pressed under 350 MPa before sintering. The pellets were pressureless sintered in a resistanceheated graphite furnace under a flowing argon atmosphere (∼1 atm) at 1950 ◦C for 60 min. The bulk density was measured with the Archimedes method. The relative density was calculated dividing the bulk density by the theoretical density that was evaluated with the rule of mixture on the basis of the starting compositions. The dense samples were examined using X-ray diffraction (Siemens D500, Germany) to identify crystalline phases. The microstructures were polished with diamond paste to  $0.25 \mu m$  and were analysed with scanning electron microscopy (SEM, Cambridge S360) and energy dispersive spectroscopy (EDS, INCA Energy 300, Oxford instruments, UK). The main properties of the processed materials are summarized in Table 1. After the arc-jet tests, the ceramic models



Fig. 2. Experimental setup. (a) Before the test and (b) during the test.

<span id="page-2-0"></span>were further analysed by SEM–EDS on surface and cross section.

# *2.2. Plasma torch tests*

Samples with a hemispheric shape ([Fig. 1\)](#page-1-0) were machined through diamond-loaded tools and then exposed to sustained enthalpy flows using the arc-jet facility equipped with a 80 kW plasma torch that operates in inert gas (He,  $N_2$ , Ar and their mixtures) at mass flow rates up to 5 g/s. The specimens were located at a distance of 6 cm from the exit torch. ([Fig. 2a](#page-1-0) and b) The HfB<sub>2</sub>-based model was tested with an initial average specific total enthalpy of about 20 MJ/kg, that was gradually increased tuning the arc current up to 26 MJ/kg and then maintained for approximately 30 s. HfC-based models were exposed to hot streams at two different conditions. The first test denoted as HC5-I, was conducted setting the initial average specific total enthalpy at 20 MJ/kg for about 40 s; then a value of 22 MJ/kg was achieved and maintained for about 60 s. During the second test, denoted as HC5-II, the specific total enthalpy was set at 22 MJ/kg during the first 90 s, 24 MJ/kg during the following 60 s and 26 MJ/kg during the last 240 s.

During the experiments, infrared and optical windows in the test chamber allowed visual inspection and diagnostic analyses. An automatic control system monitored the main parameters of the apparatus (voltage and current of the arc heater, water cooling temperature and mass flow rate). In particular, the specific total enthalpy (*H*) was evaluated through an energy balance between the energy supplied to the gas by the arc heater and the energy transferred to the cooling system (measured by the water temperature jump between inlet and outlet). The output data, processed via a dedicated software, allowed the evaluation of the surface temperature profile versus exposure time of the mode. Due to the extremely high thermal loading upon the ceramic models, surface chemical reactions like oxidation can be responsible for changes in the material's emissivity. To overcome this problem, the experiments' measurements were carried out with a radiation ratio pyrometer (Infratherm ISQ5, Impac Electronic Gmbh, Germany) which operates both in two colour and in the single colour function. In the two colour mode the instrument makes use of the ratio of two spectral radiances, measured at different wavelengths  $(0.9-1.05 \,\mu\text{m})$ , to evaluate the true temperature.

MoSi<sub>2</sub>

This overcomes the problem of the emissivity knowledge since it is supposed to be the same at both wavelengths. Once the temperature was measured with the ratio pyrometer, its value was input to evaluate the spectral emissivity using the single colour function ( $\lambda = 0.9 \,\mu$ m). In combination with the pyrometer, an infrared thermo-camera (Thermacam SC 3000, FLIR Systems, USA) was used to measure the surface temperature distributions and the spectral emissivity in the long wave range of the thermograph  $(\lambda = 9 \,\mu\text{m})$ .

# *2.3. Numerical simulation of the plasma torch flow*

To assess if the environment generated by the plasma torch at atmospheric pressure was able to reproduce heat fluxes, temperatures and reactive environments typical of atmospheric re-entry conditions, a systematic numerical analysis was carried out. The computations were carried out solving the full Navier–Stokes equations for a turbulent multi-reacting gas mixture with six chemical species (Ar, O, O<sub>2</sub>, NO, N and  $N_2$ ) in chemical nonequilibrium. Each species of the mixture was assumed to behave as a thermally perfect gas, where translational–rotational and vibrational–electronic degrees of freedom were characterised by different temperatures. Vibrational–translational energy exchanges were modelled according to the Landau–Teller model, while the vibrational relaxation time was derived from the Millikan–White formula, with Park correction for high temperatures.<sup>[18](#page-8-0)</sup> Chemical and vibrational non-equilibrium was taken into account using the Park model.<sup>[19–21](#page-8-0)</sup> The fluid dynamics equations were numerically solved in the computational domain (plasma torch and test chamber). Convective fluxes were computed according to Roe's Flux Difference Splitting scheme. Integration of the equations was implicit in time performed, until steady state was achieved, solving the linearised system of equation by the multigrid technique.

# **3. Results and discussion**

#### *3.1. Microstructural features of the as-sintered samples*

#### *3.1.1. HfB2-based composite*

 $MoSi<sub>2</sub>$ 

According to X-ray diffraction analysis (not shown), the sintered specimen was constituted by hexagonal  $HfB<sub>2</sub>$  and



 $HfO<sub>2</sub>$ 

 $(h)$ 

Fig. 3. Polished surfaces of (a)  $HfB_2$ - and (b) HfC-based composites.

<span id="page-3-0"></span>

Fig. 4. Computed mass fractions of the different species along the torch axis and specific total enthalpy contour.  $H = 26$  MJ/kg.

tetragonal  $MoSi<sub>2</sub>$  and monoclinic  $HfO<sub>2</sub>$ .<sup>[14](#page-8-0)</sup> The relative density was 98%, as reported in [Table 1. T](#page-1-0)he polished section ([Fig. 3a\)](#page-2-0) showed a regular microstructure, with little residual porosity. HfB2 grains had a rounded shape with mean grain size estimated by image analysis of about 1.5  $\mu$ m, while the MoSi<sub>2</sub> phase had an irregular morphology with very low dihedral angles (20–30◦). This peculiar characteristic indicates that MoSi<sub>2</sub> was very ductile at the sintering temperature and was accommodated between the voids left by the HfB2 skeleton. Further details are reported elsewhere.[14](#page-8-0)

#### *3.1.2. HfC-based composite*

Cubic HfC and tetragonal  $MoSi<sub>2</sub>$  were the crystalline phases detected after sintering. The final density was 98% ([Table 1\).](#page-1-0) The typical microstructure of the HfC-based composite is shown in [Fig. 3b.](#page-2-0) Very few porosity was detected by SEM inspections on

Table 2 Test conditions this material. HfC grains had a squared shape and displayed a bright colour, while the MoSi<sub>2</sub> phase had an irregular shape and was recognizable as a darker contrast phase. The mean grain size of the carbide grains was about  $3 \mu m$  ([Table 1\).](#page-1-0) Further details are reported in an earlier work.[15](#page-8-0)

# *3.2. Plasma flow characterisation*

The results herein presented refer to a 75% argon–25% nitrogen mixture plasma jet with mass flow rate of 1.45 g/s, for an average specific total enthalpy of the flow at the torch exit varying from 20 MJ/kg to 28 MJ/kg, at atmospheric pressure. At the exit of the torch the plasma containing argon, nitrogen and atomic nitrogen expands through a nozzle (5 mm in diameter), comes into contact with the surrounding air at ambient conditions, so that oxygen in the atmosphere dissociates and a reacting mixture composed of Ar,  $O_2$ ,  $N_2$ , NO, O and N is formed. Fig. 4 shows the results of computations performed for the case of an average specific total enthalpy of 26 MJ/kg. According to the calculations, the average specific total enthalpy in proximity of the specimen reduces drastically to about 6–8 MJ/kg (Fig. 4). Table 2 summarizes the test conditions and the computed flow characteristics for the different experiments. Correspondingly, [Fig. 5a](#page-4-0) and b shows the increase of the surface temperature as a function of the exposition time for  $HfB_2$  and  $HfC$  samples, respectively. It must be mentioned that the temperature reached at the sample surface depends on the ability of the material to reject the heat by radiation, i.e. on its emissivity ( $\varepsilon = 1$  for an ideal black body,  $\varepsilon$  < 1 for a real material surface). The higher is the emissivity, the greater is the emitted radiation. However, the temperature of the sample also depends on its thermal conductivity, since a high thermal conductivity allows heat to be conducted from the leading edge to colder zones and from there to be radiated away. The maximum temperature reached on the surface of HfB<sub>2</sub> sample was 1950 °C, [Fig. 5a.](#page-4-0) The corresponding stagnation point heat flux, computed by numerical simulation was in the range  $5-8$  MW/m<sup>2</sup>. A value of about 0.9 was estimated



<span id="page-4-0"></span>

Fig. 5. Temperature profiles vs. time during arc-jet testing of the (a) HB5 and (b) HC5-I and HC5-II models corresponding to specific total enthalpies during the test.

for the emissivity at the highest temperature. During the first test on HfC model (HC5-I), the surface temperature achieved by the sample was  $2050\,^{\circ}\text{C}$  (Fig. 5b). During the second test (HC5-II) the surface temperature reached the value of  $2400\degree$ C that was maintained for about 4 min. The corresponding computed surface heat flux was of the order of  $10 \text{ MW/m}^2$ . In both tests a value of about 0.7 was measured for the emissivity at the highest temperatures. The emissivity values of the samples, tested at different conditions, were found to be independent of the test temperature and conditions. These values were similar to those found for other  $ZrB_2-SiC$  ultra-high temperature ceramics tested in similar conditions.[17,22](#page-8-0)

Both the boride and carbide materials showed an excellent stability during the tests, despite the unavoidable microstructural changes occurring on their surface, as described below.

# *3.3. Microstructural modifications induced by high enthalpy plasma*

#### *3.3.1. HfB2 sample*

Due to high thermal loading and the presence of oxidising species, oxidation of the constituent phases occurred. The sample surface was covered by a compact silica-based scale (15–20  $\mu$ m thick), which embedded HfO<sub>2</sub> crystals ([Fig. 6a\)](#page-5-0). The composition of the outer glassy layer was investigated by means of EDS analysis. The silica-based scale contained several impurities, including Hf, Al and boron. Bubble formation is due to evolution of gaseous products. The shear forces associated to the hot stream enhanced the bursting of bubbles. The analysis of the cross section ([Fig. 6c\)](#page-5-0) revealed that the scale was a multilayered oxide (~70 μm thick) well adherent to the bulk, implying that no micro/macrospallation phenomena occurred. Underneath the surface silica oxide (about  $10 \mu m$ ), the scale was mainly constituted by large  $HfO<sub>2</sub>$  grains and silica with composition close to  $SiO<sub>2</sub>$  [\(Fig. 6e\)](#page-5-0). Occasionally Al impurities were detected inside the glassy phase. In the innermost layer formation of molybdenum oxides, silicon oxides and hafnium oxide was observed. Nitrogen impurities were also detected in the composition of the molybdenum-based oxide. The chemistry of the experiment carried out on  $HfB_2$ -based material was governed by the oxidation reaction of the two constituent phases. Hafnium diboride oxidises according to  $2^3$ :

$$
HfB_2(s) + (5/2)O_2(g) = HfO_2(s) + B_2O_3(l)
$$
 (1)

$$
B_2O_3(l) = B_2O_3(g)
$$
 (2)

Hafnia is a very stable phase in oxidising atmosphere above 2000 °C. It has a melting point of 2900 °C and relatively low vapour pressure.<sup>24</sup> Boron oxide has a low melting point and high vapour pressure, therefore at  $T > 1100$  °C, it starts to evaporate, according to reaction (2). On the other hand, at temperatures >  $1000\degree C$ , MoSi<sub>2</sub> is known to form a stable silica layer according to  $2^5$ :

$$
MoSi2 + (7/2)O2(g) = MoO3(g) + 2SiO2
$$
 (3)

The nature of the oxide observed suggests that during the plasma torch tests the silica production according to reaction (3) was fast enough to protect the material. The formation of bubble was produced by the escape of gaseous by-products, such as  $MoO<sub>3</sub>, B<sub>2</sub>O<sub>3</sub>$  and other highly volatile boron suboxides. However, the presence of B impurities in the outer glassy layer suggests that due to the short exposure the release of boron gaseous products was not complete. The role of dissociated oxygen, which is the primary oxidant agent in a re-entry environment, is still matter of debate. Substituting atomic oxygen to molecular oxygen, thermodynamic calculations (HSC Chemistry for Windows 5, Outokompu Research, OY, Pori, Finland) indicate that reactions (2) and (3) are even more favoured (the absolute value of free Gibbs energy increases by a factor comprised between 1.4 and 1.7). Preliminary results reported in the literature on  $\text{ZrB}_2-\text{SiC}$  composites subjected to re-entry simulations confirmed that oxidation by atomic oxygen proceeds more rapidly than by molecular oxygen.[4](#page-7-0)

Post-tests analyses unfortunately do not help to clarify these aspects. The morphology of the cross section, despite the drastic differences in the experimental conditions, is very similar to that of  $HfB_2-MoSi_2$  composites oxidised under conventional conditions, such as static air, longer holding times (15 min) and temperatures of 1400–1650 °C. Recent studies<sup>[11](#page-8-0)</sup> confirm that the response of  $ZrB_2-SiC$  and  $HfB_2-SiC$  ceramics to arc heater

<span id="page-5-0"></span>

Fig. 6. (a) HfB2 sample after arc-jet test, (b) surface, (c) cross section, (d) enlarged view of layers I, II and (e) Enlarged view of layer III.

testing appear to be similar to conventional oxidation studies. Finally, it is worth noting that differently from  $\text{ZrB}_2-\text{SiC}$  and HfB2–SiC composites no depletion layer was observed in the scale, i.e. no active oxidation of MoSi<sub>2</sub> occurred for this system.

# *3.3.2. HfC samples*

After the HFC-I tests no significant variation of the sample size and shape was observed, implying that the extent of ablation was very low. The surface turned from dark grey to a whitish colour [\(Fig. 7a](#page-6-0) and b) and was constituted by a visibly cracked hafnium oxide scale, with no glassy phase. The specimen section [\(Fig. 7c](#page-6-0)) displayed the formation of a multilayered scale, with thickness of about 90  $\mu$ m, well adherent to the unreacted bulk. Underneath the surface, the outermost portion of the scale was constituted by hafnium oxide and a silica-based glassy phase, which partially filled the porosity ([Fig. 7d](#page-6-0)). The intermediate layer contained a fine porosity and was constituted by hafnia and isolated pockets of molybdenum oxide [\(Fig. 7e\)](#page-6-0). The innermost layer contained partially oxidised HfC and SiC, and residual Mo–Si species ([Fig. 7f\)](#page-6-0). No porosity was found in this region. Similar features were displayed by the model tested at 2400 ◦C (test HfC-II), even if the oxide scale ( $\sim$ 300  $\mu$ m thick) was found to be partially detached from the unreacted bulk.

Despite the presence of MoSi<sub>2</sub> as a  $SiO<sub>2</sub>$ -forming phase, no evidence of continuous glassy layer was found on the surface of HfC, in contrast with results obtained for  $HfB<sub>2</sub>$ . This finding cannot be totally imputed to the extreme conditions of the present experiments. The two compositions, HB5 and HC5, were in fact oxidised at the same conditions, in static air (i.e. in absence of significant ablation) at  $1650^{\circ}$ C, in a conventional furnace. Even under these milder conditions,  $HfB<sub>2</sub>$  displayed the formation of a compact layer of silica, while the surface of HfC was mainly covered by an  $HfO<sub>2</sub>$  scale with few discontinuous pockets of silica. This experiment suggests that the presence of a carbide matrix rather than a boride apparently hindered the formation of a stable silica layer. It can be hypothesized that CO species deriving from oxidation of HfC interacted with  $M_0Si<sub>2</sub>$  causing release of volatile SiO.

In the cross section, the layered structure of the oxide ([Fig. 7\)](#page-6-0) resemble the results presented in the literature on the oxidation behaviour of monolithic  $HfC^{24,26}$  $HfC^{24,26}$  $HfC^{24,26}$  These studies reported the formation of a layered scale, which comprised a porous outer layer, a dense interlayer and an oxycarbide layer, HfO<sub>2−*x*</sub>C<sub>*y*</sub>.<sup>[24,26](#page-8-0)</sup> The addition of  $MoSi<sub>2</sub>$  in the composite of the present work resulted in formation of silica which partially filled the inner porosity of the  $HfO<sub>2</sub>$  scale, a feature which should improve its oxidation resistance.

In the innermost layer,  $HfO_xC_y$  species and  $SiO_xC_y$  species were observed, the latter at the interface between MoSi<sub>2</sub> and HfC. The formation of  $SiO<sub>x</sub>C<sub>y</sub>$  species could to be related again to interaction of CO species with  $M_0Si_2$  at the very low oxygen partial pressure existing under the scale. The oxidation mechanisms of this composite still have to be completely understood and will be object of further investigation.

As for the diboride system, the role of atomic oxygen is unclear. In the literature no studies can be found on this topic. The microstructural features of the HfC-based material after testing in the arc-jet facility are indeed very similar to those of samples oxidised in conventional furnaces.

<span id="page-6-0"></span>

Fig. 7. (a) HfC sample after arc-jet test, (b) surface, (c) cross section. Enlarged view of (d) layer I (SE imaging), (e) layer II (BSE imaging), (f) layer III (BSE imaging).

# **4. Numerical rebuilding for surface catalytic efficiency evaluation**

Numerical computations were carried out using the model described in Section [2.3,](#page-2-0) under different assumptions about the catalytic properties of the specimen surface with reference to catalytic efficiency of atomic nitrogen and oxygen. The simulations refer to the four plasma torch test conditions reported in [Table 2.](#page-3-0) Based on the computed heat flux distributions, a thermal analysis was carried out to evaluate the catalytic efficiency value  $\gamma$  needed to fit the experimental temperature data. The catalytic efficiency is defined as the ratio of the number of dissociating atoms that recombine at the wall to the total number of the colliding atoms with the wall. For a non catalytic wall this value is 0, while for a fully catalytic wall this is 1. Fig. 8 shows the steady state results computed under the two assumptions of fully catalytic and non catalytic wall and the experimental data obtained with the pyrometer for both material samples. For thermal analysis, input values of specific heat and thermal conductivity are necessary. As a first approximation, the values of monolithic HfC and HfB<sub>2</sub> were considered for the composites, as follows:  $200 \text{ J/(kg} \degree C)$  and  $22 \text{ W/m} \degree C$  for specific heat and



Fig. 8. Experimental results and numerical solutions corresponding to the different assumption of fully catalytic (FC) and non catalytic (NC)wall for (a) HfB2 and (b) HfC.

<span id="page-7-0"></span>

Fig. 9. Numerical evaluation of coefficient of catalytic recombination for HfC- $5\%$  MoSi<sub>2</sub> as a function of the temperature.

thermal conductivity of HfC,<sup>[27](#page-8-0)</sup> 300 J/(kg  $\degree$ C) and 80 W/m  $\degree$ C for specific heat and thermal conductivity of  $HfB_2$ <sup>[28](#page-8-0)</sup>. The data displayed in [Fig. 8a](#page-6-0) and b, at the same plasma torch conditions, i.e. the same free stream conditions, highlight that the heating behaviour of the two materials was different. The experimental results of HfB<sub>2</sub> sample matched well the numerical values corresponding to the non catalytic wall condition. This points out that the material herein tested exhibits a non catalytic behaviour at very high temperatures. This behaviour can be explained by the formation of a silica surface layer [\(Fig. 6\)](#page-5-0) which is known to possess very low catalytic recombination behaviour.<sup>[29,30](#page-8-0)</sup> The presence of such a surface layer also justifies the high values of the surface emissivity, according to the literature data.

The experimental results for HfC suggested a partially catalytic behaviour. In order to identify a dependence of surface catalytic efficiency with temperature, different values were considered for each test conditions. Fig. 9 shows the results of the computations. At  $1800\,^{\circ}\text{C}$  HfC exhibited a non catalytic behaviour. Increasing the temperature the catalytic efficiency increased up to a value of  $2 \times 10^{-3}$  at 2400 °C [\(Fig. 4a](#page-3-0) and b) which is relatively low with respect to the fully catalytic wall condition  $(y=1)$  and of the same order of other low catalytic materials such as those of the Space Shuttle tiles. $31$ 

It should be pointed out that the present tests have been carried out at atmospheric pressure conditions. Experimental and theoretical works on the catalytic activity of silica-based mate-rials under simulated re-entry conditions<sup>[32,33](#page-8-0)</sup> showed that at constant temperature the catalytic atomic recombination coefficients are decreasing functions of the pressure. Therefore the catalytic properties of the material, in respect to the recombination of oxygen atoms, may be larger at lower pressures, as found for instance in arc-jet experiments with  $\text{ZrB}_2/\text{SiC}$  and  $\text{HfB}_2/\text{SiC}$ ceramic materials.<sup>2</sup>

#### **5. Summary and future work**

Two different ultra-high temperature ceramics,  $HfB_2 + 5\%$  $M_0$ Si<sub>2</sub> and HfC + 5% MoSi<sub>2</sub> were produced by pressureless sintering. Machined hemispherical models were exposed to ground simulated atmospheric re-entry conditions using arc-jet testing, with an average specific total enthalpy of the flow around the body of the order of 5–10 MJ/kg and at atmospheric pressure.

The  $HfB_2 + 5\%$  MoSi<sub>2</sub> model surface reached a peak value of 1950 ◦C for *H* approaching 8 MJ/kg. SEM–EDS analysis of the cross section after exposure showed the formation of a compact silica oxide (about 15  $\mu$ m) which sealed the underlying HfO<sub>2</sub> scale.

The HfC  $+5\%$  MoSi<sub>2</sub> model surface reached peak values of 2100  $\degree$ C and 2400  $\degree$ C. Cross section analysis showed a layered structure, constituted of an outer layer of porous  $HfO<sub>2</sub>$  and an inner layer mainly constituted of  $HfO<sub>2</sub>$  and silica.

Numerical calculations, which simulated the chemical nonequilibrium flow around the hemispheric model correlated well with the experimental results assuming a very low catalytic surface behaviour for  $HfB_2$  and a catalytic behaviour increasing with temperature for HfC.

Although more testing is necessary to improve our understanding of the oxidation mechanisms under extreme conditions, the composites presently tested showed an excellent resistance to high enthalpy hot flows. This stability at temperature around  $2000\degree$ C opens up new developments in several fields of application, including nuclear applications and industries where extreme conditions are involved. Moreover, the possibility to produce near net shape components through a conventional sintering technique represents a technological advantage in comparison with materials in need of pressure-assisted techniques and expensive post-sintering machining.

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